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### **Radiationless Deactivation of Singlet Excited States of Complex Molecules in the Gas Phase**

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RADIATIONLESS DEACTIVATION OF SINGLET EXCITED STATES  
OF COMPLEX MOLECULES IN THE GAS PHASE

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ABSTRACT

The ways of degradation of electronic excitation energy of singlet states of complex molecules in the gas phase with various character of dependence of the quantum yield and lifetime of fluorescence on the excess vibrational energy of excited molecules,  $\Delta E_{\text{vib}}^*$ , are discussed. From the analysis of literature data on phthalimide and anthracene derivatives and perylene, and also from the results of investigations conducted in the present study for 9,10-dimethylantracene and 9,10-diphenylantracene it follows that in the majority of cases radiationless degradation of electronic excitation energy is realized both by means of intersystem crossing and internal conversion. For all studied molecules, with the exception of anthracene, an increase in intersystem crossing with increasing  $\Delta E_{\text{vib}}^*$  of excited singlet state is observed, while the internal conversion probability may either increase or decrease. In particular, an

increase in the quantum yield and lifetime of fluorescence of 9,10-dimethylantracene with increasing  $\Delta E_{\text{vib}}^*$  proved to be due to a decrease in the internal conversion probability.

An increase in the excess vibrational energy of singlet excited states of the majority of polyatomic organic compounds in the gas phase is known to be accompanied with an increase in the probability of radiationless transitions<sup>1</sup>. As a result, the quantum yield,  $\phi$ , and lifetime,  $\tau$ , of fluorescence of free complex molecules usually decrease with increasing energy of the exciting radiation quantum,  $h\nu_{\text{exc}}$ , and vapour temperature,  $T$ .

Radiationless degradation of electronic energy of the first excited singlet  $S_1$  state can be realized in two ways: by means of intersystem crossing ( $S_1 \rightarrow T_n$  transition), and internal conversion ( $S_1 \rightarrow S_0$  transition). The role of each of the channels in the radiationless dissipation of electronic excitation energy of free polyatomic molecules, however, remained uninvestigated until recently. In our studies<sup>2,3</sup> it has been shown for a series of phthalimide derivatives in the gas phase that the  $S_1$  state energy degradation is realized through both channels. Moreover, with increasing excess vibrational energy of the excited molecules,  $\Delta E_{\text{vib}}^*$ , the intersystem crossing probability usually increases more rapidly, thus causing a sharp decrease in  $\phi$  and  $\tau$ .

However, there are some compounds whose  $\phi$  remains practically constant with increasing  $\Delta E_{\text{vib}}^*$  (anthracene,

perylene), or even increases (9,10-dimethylantracene)<sup>1</sup>. For the anthracene molecule in the gas phase intersystem crossing proved to be the main channel of radiationless deactivation of the  $S_1$  state, its probability being independent of  $\Delta E_{vib}^*$ <sup>3</sup>. In the case of perylene molecules in the gas phase, internal conversion effectively competes with intersystem crossing in the course of radiationless dissipation of the electronic excitation energy. The probabilities of these processes change so that their sum remains practically constant<sup>4</sup>.

The principle aim of the present study was to establish the ways of the degradation of electronic energy of the  $S_1$  states of complex molecules in the case when the fluorescence quantum yield and lifetime increase with increasing  $\Delta E_{vib}^*$ . Vapours of 9,10-dimethylantracene (9,10-DMA) and 9,10-diphenylantracene (9,10-DPA) were studied. The first compound is characterized by an increase of the values of  $\gamma$  and  $\tau$  with increasing excess vibrational energy of the excited molecules,  $\Delta E_{vib}^*$ , while the second compound, by their decrease. The triplet-triplet absorption spectra and quantum yields of triplet state formation of the vapours of these compounds were measured in order to determine the probabilities of intersystem crossing,  $d_{ST}$ , and internal conversion,  $d_{SS}$ .

The molecules under study were excited by the second harmonic radiation of a giant-pulse ruby laser. Triplet-triplet absorption was recorded 100-200 nsec after the excitation of molecules. Photoelectric method

was used to determine the intensity of the radiation of the spectroscopic flash after passing through an unexcited sample,  $I_{\text{unexc}}$ , and through the excited one,  $I_{\text{exc}}$ . The experimentally measured ratio,  $I_u/I_e$ , was used to determine the excited sample optical density:

$$D^T = 1 - \frac{I_u}{I_e} = (\alpha^T - \alpha^S)n_T l,$$

where  $n_T$  is the triplet state population;  $\alpha^S$  and  $\alpha^T$ , molecular coefficients of singlet and triplet absorption, respectively;  $l$ , optical path length in the sample (14 cm). The excess vibrational energy was varied by means of introduction into the vapour of the foreign gas - pentane - at different pressures. It is known<sup>5</sup> that the interaction of foreign gases with molecules in electronic excited states results in vibrational energy exchange which leads to the change of radiationless transition probabilities. The studied molecules, excited by the ruby laser second harmonic radiation into the  $S_1$  state, may be considered isolated under the vapour pressures used, and have the vibrational energy excess of  $2600 \text{ cm}^{-1}$ . As a result of introducing of the foreign gas (400 mm Hg) into the vapour, the practically equilibrium distribution of molecules over the  $S_1$  state vibrational levels is realized. Figure 1 displays spectra of triplet-triplet absorption of 9,10-DMA and 9,10-DPA in pure vapour and in the presence of the foreign gas.

To determine  $n_T$  and  $\alpha^T$ , we used here as well as in our previous works the method suggested in Ref.<sup>6</sup>. The absolute values of the triplet state formation quantum

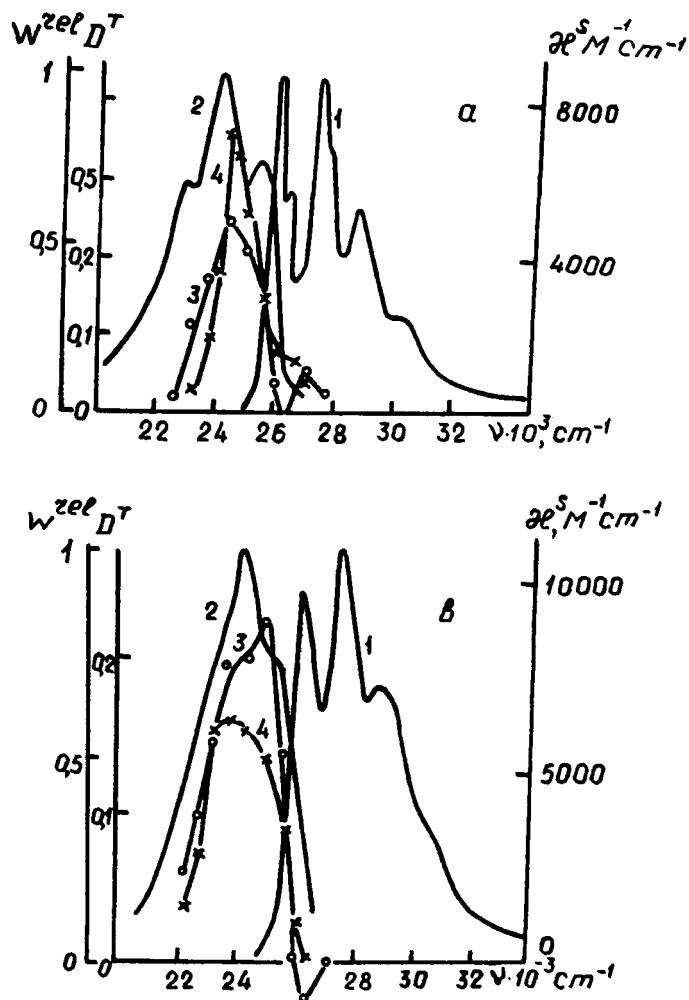


Fig.1. Spectra of singlet absorption (1), fluorescence (2) and optical density,  $D^T$ , (3, 4 - at different pressures of pentane,  $P_p$ ):  
 a - 9,10-dimethylantracene vapours ( $P = 0.23$  mm Hg,  $T = 448$  K), 3 -  $P_p = 0$ ;  
 4 -  $P_p = 360$  mm Hg; b - 9,10-diphenylantracene vapours ( $P = 0.3$  mm Hg,  
 $T = 525$  K), 3 -  $P_p = 0$ , 4 -  $P_p = 400$  mm Hg.

yield,  $\eta$ , were found from the ratio of the concentration of triplet molecules to that of molecules which absorbed the exciting radiation. In these experiments nearly a half of the incident exciting radiation energy was absorbed in the cuvette. The exciting radiation power density was maintained practically constant along the entire cuvette length by means of decreasing of the cross-section area of the light beam from the beginning of the cuvette towards its end with the help of a focusing lens. No more than 15% of molecules present in the excited volume were converted into the triplet state. The error of determination of the absolute values of  $\eta$  did not exceed 40%, errors in relative measurements of  $\eta$  were less than 15%. Using the values of  $\eta$  obtained and those of  $\gamma$  and  $\tau$  from the literature data, we estimate the probabilities of intersystem crossing,  $d_{ST}$ , and internal conversion,  $d_{SS}$ , from the  $S_1$  states of the molecules under study, at various excesses of vibrational energy (see Table 1).

In the case of 9,10-dimethylanthracene, radiationless degradation of electronic energy proceeds through both channels. The internal conversion yield in the  $S_1$  state decreases with increasing excess vibrational energy (within the studied limits) from 0.52 to 0.12, the internal conversion probability,  $d_{SS}$ , is reduced threefold, while the intersystem crossing probability,  $d_{ST}$ , increases approximately by a factor of 1.5. Thus an increase of  $\gamma$  and  $\tau$  with increasing  $\Delta E^*_{vib}$  in this case is due to a more rapid decrease of  $d_{SS}$  as compared to the increase of  $d_{ST}$ .

TABLE 1

Intersystem Crossing and Internal Conversion Probabilities

Compound	Vapour tempe- rature, K	T-T absorption coefficient, $\epsilon^T, M^{-1}cm^{-1}$	T-state quantum yield, $\eta$	Fluo- rescen- ce quantum yield, $\phi [7]$	S <sub>1</sub> state life- time, $\tau$ ns [7]	$d_{ST} \cdot 10^{-8}$ s <sup>-1</sup>	$d_{SS} \cdot 10^{-6}$ s <sup>-1</sup>
9,10-di- methyl- anthra- cene	448	11900 $\pm$ 1700 ( $\nu \approx 24390 \text{ cm}^{-1}$ ) nonrelaxating molecules	0.67	0.14	3.5	1.9	0.5
		30100 $\pm$ 4500 ( $\nu \approx 24390 \text{ cm}^{-1}$ ) relaxating molecules (360 mm Hg pentane)	0.37	0.11	3.0	1.3	1.7
		12400 $\pm$ 1800 ( $\nu \approx 25000 \text{ cm}^{-1}$ ) nonrelaxating molecules	0.34	0.76	3.3	1.0	-
		11000 $\pm$ 1600 ( $\nu \approx 23800 \text{ cm}^{-1}$ ) relaxating molecules (400 mm Hg pentane)	0.26	0.8	3.3	0.8	-
9,10-di- phenyl- anthra- cene	525	12400 $\pm$ 1800 ( $\nu \approx 25000 \text{ cm}^{-1}$ ) nonrelaxating molecules	0.34	0.76	3.3	1.0	-
		11000 $\pm$ 1600 ( $\nu \approx 23800 \text{ cm}^{-1}$ ) relaxating molecules (400 mm Hg pentane)	0.26	0.8	3.3	0.8	-



It can be seen from Fig.1a that the optical density,  $D^T$ , of 9,10-DMA vapours in the vicinity of the triplet-triplet absorption band maximum in the presence of pentane exceeds  $D^T$  of rarefied vapours, although the triplet state formation yield is reduced on a decrease of excess vibrational energy of singlet-excited molecules caused by the foreign gas addition (Table 1). This is due to an increase of the triplet-triplet absorption coefficient resulting from vibrational relaxation of complex molecules during interaction with the foreign gas molecules. The increase of triplet-triplet absorption coefficient near the band maximum successfully competes with the decrease of triplet molecules concentration,  $n_T$ , due to a decrease of the probability  $d_{ST}$  upon addition of pentane to the 9,10-DMA vapours. This can be seen from the time dependencies of  $D^T$  (Fig.2). In the case of rarefied vapours, within the investigated time interval, the intermolecular vibrational relaxation in the triplet state does not occur, and  $D^T$  remains constant (curve 3). Contrary to this, with the pentane pressure being 360 mm Hg, the relaxation is already completed (curve 2). In the presence of pentane at 13 mm Hg, the molecules of 9,10-DMA during their lifetime in the excited singlet state undergo approximately one collision, and the probability  $d_{ST}$  must be less than in the case of rarefied vapours. Consequently, at the initial instant of observation,  $D^T$  (curve 1) is also less than for rarefied gases. With time,  $D^T$  increases due to an increase in triplet-triplet absorption coefficient during intermolecular relaxation in the triplet state,

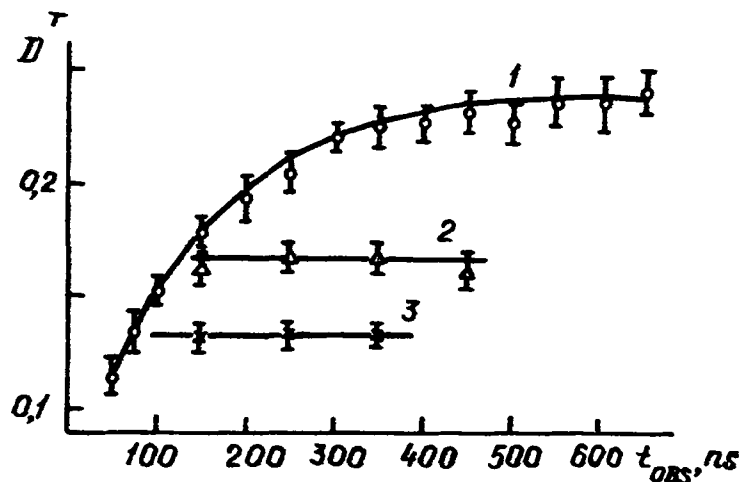


Fig. 2. Dependence of triplet-triplet absorption of 9,10-dimethylantracene vapours at the frequency  $\nu = 24390 \text{ cm}^{-1}$  on time at different pressures of pentane (1 - 13, 2 - 360, 3 - 0 mm Hg)

and becomes larger than  $D^T$  at 360 mm Hg of pentane, which is attributable to a considerably lower intersystem crossing probability at high pressure of pentane.

The sum of quantum yields of triplet state formation and fluorescence of 9,10-diphenylantracene vapours is equal to unity within experimental error limits, which means that the radiationless deactivation of the  $S_1$  state of this compound proceeds by intersystem crossing and its probability increases with an increase in excess vibrational energy, thus causing the decrease of the fluorescence quantum yield.

From the analysis of the data obtained by us for different classes of compounds in the gas phase (phtha-

limide and anthracene derivatives, perylene) it follows that in the majority of cases radiationless degradation of electronic excitation energy is realized through both channels. For all molecules studied, with the exception of anthracene, an increase in intersystem crossing probability with increasing excess vibrational energy of the excited state is observed, while the probability of internal conversion may either increase or decrease, being practically negligible for some molecules. This does not agree with the published calculation data<sup>3,9</sup>, according to which the internal conversion probability must increase with increasing excess vibrational energy of singlet-excited molecules.

In the solutions of the studied aromatic hydrocarbons - perylene, anthracene, 9,10-dimethylantracene and 9,10-diphenylantracene - the channel of electronic excitation radiationless degradation is the intersystem crossing<sup>10</sup>. This result agrees with the empirical regularity established for solutions of aromatic hydrocarbons<sup>11,12</sup>, according to which the internal conversion probability is reduced with an increase in the energy interval between the combining singlet electronic states. Since the value of energy gap between the  $S_0$  and  $S_1$  levels of the above-mentioned compounds exceeds  $20000\text{ cm}^{-1}$ , the internal conversion probability must not be more than  $10^6\text{ s}^{-1}$  on the order of magnitude, which is much less than the sum of probabilities of radiation transition and intersystem crossing,  $10^3\text{ s}^{-1}$ .

As we have shown earlier, in the case of perylene and 9,10-dimethylantracene vapours, unlike the conden-

sed phase, internal conversion is an effective channel of radiationless energy degradation. A possible reason for this behaviour of free molecules, the importance of which was underlined by Terenin<sup>13</sup>, is the enhancement in the gas phase of torsional and deformational vibrations causing a greater molecular skeleton mobility than in the condensed phase.

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